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Radio spectral properties of star-forming galaxies in the MIGHTEE-COSMOS field and their impact on the far-infrared-radio correlation

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ABSTRACT

We study the radio spectral properties of 2094 star-forming galaxies (SFGs) by combining our early science data from the MeerKAT International GHz Tiered Extragalactic Exploration (MIGHTEE) survey with VLA, GMRT radio data, and rich ancillary data in the COSMOS field. These SFGs are selected at VLA 3 GHz, and their flux densities from MeerKAT 1.3 GHz and GMRT 325 MHz imaging data are extracted using the ‘superdeblending’ technique. The median radio spectral index is $\alpha_{1.3\text{ GHz}}^{3\text{ GHz}} = -0.80 \pm 0.01$ without significant variation across the rest-frame frequencies ~ 1.3 –10 GHz, indicating radio spectra dominated by synchrotron radiation. On average, the radio spectrum at observer-frame 1.3–3 GHz slightly steepens with increasing stellar mass with a linear fitted slope of $\beta = -0.08 \pm 0.01$, which could be explained by age-related synchrotron losses. Due to the sensitivity of GMRT 325 MHz data, we apply a further flux density cut at 3 GHz ($S_{3\text{ GHz}} \geq 50 \mu\text{Jy}$) and obtain a sample of 166 SFGs with measured flux densities at 325 MHz, 1.3 GHz, and 3 GHz. On average, the radio spectrum of SFGs flattens at low frequency with the median spectral indices of $\alpha_{325\text{ MHz}}^{1.3\text{ GHz}} = -0.59_{-0.03}^{+0.02}$ and $\alpha_{1.3\text{ GHz}}^{3.0\text{ GHz}} = -0.74_{-0.02}^{+0.01}$. At low frequency, our stacking analyses show that the radio spectrum also slightly steepens with increasing stellar mass. By comparing the far-infrared-radio correlations of SFGs based on different radio spectral indices, we find that adopting $\alpha_{1.3\text{ GHz}}^{3\text{ GHz}}$ for k -corrections will significantly underestimate the infrared-to-radio luminosity ratio (q_{IR}) for >17 per cent of the SFGs with measured flux density at the three radio frequencies in our sample, because their radio spectra are significantly flatter at low frequency (0.33–1.3 GHz).

Key words: methods: observational – galaxies: evolution – galaxies: formation – radio continuum: galaxies.

1 INTRODUCTION

The radio emission from star-forming galaxies (SFGs) consists of non-thermal synchrotron radiation from the cosmic ray (CR) electrons spiralling in the interstellar magnetic field, and thermal free–free emission produced by the electrostatic interactions between charged particles, mostly free electrons and ions, in H II regions. In SFGs, the CR electrons are primarily accelerated by shocks associated with supernovae remnants. In addition, the radio continuum emission at $\nu \lesssim 30$ GHz is unaffected by dust attenuation. These characteristics make radio emission a key observable that provides information about the CR electrons, magnetic fields, photoionization rate, dust-obscured star formation, etc. in galaxies (see Condon 1992, for a review). Therefore, deep radio continuum surveys are essential for our understanding of physical processes at work in SFGs from the local to the distant Universe (e.g. Ibar et al. 2009; Murphy 2009; Beswick et al. 2015; Prandoni & Seymour 2015; Jarvis et al. 2016; Smolčić et al. 2017a; Ocran et al. 2020a).

In the last four decades, numerous studies have focused on characterizing the radio continuum spectra of SFGs to explore the relative contributions of synchrotron and free–free emission, components or structures corresponding to these two radiation mechanisms, and their correlation with emission in other wavebands (e.g. Klein, Wielebinski & Morsi 1988; Richards et al. 1998; Haarsma et al. 2000; Condon, Cotton & Broderick 2002; Seymour et al. 2008; Williams & Bower 2010; Marvil, Owen & Eilek 2015; Tabatabaei et al. 2017; Klein, Lisenfeld & Verley 2018; Gim et al. 2019). Traditionally, the radio spectrum of an SFG is assumed to be a superposition of a steep synchrotron spectrum with spectral index $\alpha = -0.8$ ($S_\nu \propto \nu^\alpha$) and a flat free–free spectrum with $\alpha = -0.1$, which is also supported by observations of some nearby SFGs with a star formation rate (SFR) $< 10 M_\odot \text{ yr}^{-1}$ (e.g. Condon 1992; Niklas, Klein & Wielebinski 1997; Tabatabaei et al. 2017). However, studies of luminous and ultraluminous infrared galaxies (ULIRGs), starburst galaxies, dwarf galaxies, or even some normal SFGs in the local Universe and at high redshift have shown that the radio continuum spectral energy distributions (SEDs) of SFGs are rarely characterized well by a single power law (e.g. Clemens et al. 2010; Williams &

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Bower 2010; Murphy et al. 2013; Marvil et al. 2015; Klein et al. 2018; Thomson et al. 2019; Tisanić et al. 2019).

A well-determined radio spectrum for SFG is critically important for studies that are based on the rest-frame radio power, especially those at high redshift, which are most sensitive to the assumed k -corrections. These include studies of the correlation observed between radio and far-infrared (FIR) emission of SFGs, which is widely adopted to calibrate the SFR measured from radio continuum emission, and the definition of the ‘radio excess’ that is used to classify active galactic nuclei (AGNs) and SFGs (e.g. Helou, Soifer & Rowan-Robinson 1985; Yun, Reddy & Condon 2001; Bell 2003; Ivison et al. 2010; Mao et al. 2011; Del Moro et al. 2013; Magnelli et al. 2015; Hindson et al. 2018; Ocran et al. 2020a,b, 2021). The FIR-radio correlation (FIRRC) has been well-established for local SFGs at GHz frequencies. However, for SFGs at high redshift, the radio emission at rest-frame ~ 1.4 GHz is shifted to low frequencies ($\nu_{\text{obs}} < 1$ GHz at $z > 0.4$). Therefore, the appropriate spectral index that is used to k -correct the observed radio flux densities at ~ 1.4 GHz for SFGs at $z > 0$ should be drawn from lower frequency ($\nu_{\text{obs}} < 1.4$ GHz) data. For SFGs, previous work demonstrated that at $\nu < 1$ GHz, the H II region becomes optically thick and the relativistic bremsstrahlung, synchrotron self-absorption, and Razin effects become more important than that at high frequency (e.g. Rybicki & Lightman 1979; Condon 1992; Lacki 2013; Chyży et al. 2018). These effects can suppress the radio emission and hence flatten the radio spectrum at low frequency, as has been observed in nearby SFGs (e.g. Condon 1992; Clemens et al. 2010; Murphy et al. 2013; Marvil et al. 2015; Kapińska et al. 2017; Chyży et al. 2018).

For the high-redshift Universe, although some recent studies also found a flatter radio spectrum at low frequency, these works only focus on the high-SFR ($\text{SFR} > 100 M_{\odot} \text{ yr}^{-1}$, Tisanić et al. 2019) or submillimeter bright galaxies (Thomson et al. 2019). In addition, due to the lack of high-quality radio data at low frequency, most of the studies based on rest-frame radio spectra of SFGs at high redshift have adopted a spectral index measured from high frequency ($\nu_{\text{obs}} = 1\text{--}10$ GHz), or even assumed a fixed spectral slope, when k -correcting the observed flux densities (e.g. Appleton et al. 2004; Ibar et al. 2010; Ivison et al. 2010; Sargent et al. 2010; Mao et al. 2011; Delhaize et al. 2017; Algera et al. 2020; Delvecchio et al. 2021). However, some recent works suggest that these assumptions will dramatically overestimate the radio flux density at low frequency and thus strongly affect the FIRRC of SFGs (e.g. Schleicher & Beck 2013; Delhaize et al. 2017; Galvin et al. 2018; Gim et al. 2019). Therefore, high-quality radio observations at both low and high frequency from the Square Kilometre Array (SKA)¹ and its precursors, such as MeerKAT and ASKAP, will be extremely important in accurately determining the radio spectra of SFGs at high redshift and then shedding light on galaxy formation and evolution through the radio window (Beswick et al. 2015; Prandoni & Seymour 2015; Jarvis et al. 2016).

In this work, we use the early science data from the MeerKAT International GHz Tiered Extragalactic Exploration (MIGHTEE) survey and take advantage of existing radio data from *Karl G. Jansky* Very large Array (VLA) and Giant Metrewave Radio Telescope (GMRT) surveys, as well as the rich ancillary data in the COSMOS field to select SFGs up to $z \sim 3$ and study their radio spectral properties. By studying the shape of radio spectra with the physical properties derived from other bands, we discuss the mechanisms that

characterize the radio spectra at low and high frequencies and their impact on the FIRRC of SFGs.

The radio and other ancillary data used in this work are described in Section 2. We present our analyses of radio spectral properties based on these data in Section 3. Our results are shown in Section 4. We discuss and summarize our results in Sections 5 and 6, respectively. Throughout this paper, we adopt the AB magnitude system (Oke 1974) and assume a flat Λ CDM cosmological model with the Hubble constant $H_0 = 67.27 \text{ km s}^{-1} \text{ Mpc}^{-1}$, matter density parameter $\Omega_m = 0.32$, and cosmological constant $\Omega_{\Lambda} = 0.68$ (Planck Collaboration XIII 2016).

2 OBSERVATIONS AND DATA

2.1 MIGHTEE 1.3 GHz survey in the COSMOS field

The MIGHTEE survey is one of the MeerKAT large survey projects.² The main purpose of MIGHTEE is to obtain deep GHz radio continuum, spectral line, and polarization observations to study the cosmic evolution of galaxies and AGNs. Details of the science goals and observation plan of the MIGHTEE survey are described in Jarvis et al. (2016). Taking advantage of existing multiwavelength data, the MIGHTEE project is surveying four well-studied extragalactic fields, i.e. COSMOS (1 deg^2), ELAIS S1 (1.6 deg^2), CDFS (8.3 deg^2), and XMM-LSS (6.7 deg^2) fields, with a median sensitivity of $2 \mu\text{Jy beam}^{-1}$ at L-band (900–1670 MHz) and $1 \mu\text{Jy beam}^{-1}$ at S-band (1750–3500 MHz) with the latter only for the CDFS and the COSMOS fields. We refer the reader to Jarvis et al. (2016) for more details.

In this work, we use the early science data of the COSMOS field (R.A.: 10:00:28.6, Decl.: +02:12:21), which were observed in 2018 and 2020 with the MeerKAT L-band receiver ($\nu_c = 1.3$ GHz). Because of the frequency dependence of the primary beam as well as the wide bandwidth, the effective frequency gradually decreases from the centre outwards. However, the MIGHTEE-COSMOS early science data were observed with a single pointing and the decrease of the effective frequency is < 0.04 GHz within the primary beam ($\sim 1 \text{ deg}$ FWHM). This variation is negligible in studying radio spectra using MeerKAT, VLA 3 GHz, and GMRT 325 MHz data. Therefore, we use $\nu_c = 1.3$ GHz as the effective frequency of MeerKAT L-band data in this work.

In total, the observations include 17.45 h on source for the central pointing in the COSMOS field. Details of the observations and data reduction will be presented in I. Heywood et al. (in preparation) but see also a summary in Delhaize et al. (2021). Here, we provide a brief overview of the final reduced early science data.

Because high-sensitivity and high-resolution cannot be achieved at the same time, the MIGHTEE continuum data are imaged twice with Briggs’ robust parameter 0.0 and -1.2 . In this work, to obtain a less-biased sample of SFGs for the study of their radio spectral properties (Section 4), we only use the high-sensitivity imaging data. The maximum-sensitivity image reaches a thermal noise of $1.7 \mu\text{Jy beam}^{-1}$ with an angular resolution of $8''.6 \times 8''.6$. In addition to the thermal noise, the contribution from the numerous faint and unresolved astronomical sources, namely classical confusion, increases the mean root-mean-square (RMS) to $4.5 \mu\text{Jy beam}^{-1}$.

¹<https://www.skatelescope.org/>

²<http://public.ska.ac.za/meerkat/meerkat-large-survey-projects>

2.2 Other radio data in the COSMOS field

2.2.1 VLA 3 GHz and 1.4 GHz data

The *Karl G. Jansky* VLA-Cosmic Evolution Survey (VLA-COSMOS) 3 GHz Large Project mapped the entire 2 deg^2 COSMOS field with a median RMS of $2.3 \mu\text{Jy beam}^{-1}$ and an angular resolution of $0''.75$ (Smolčić et al. 2017a). In total, 10 830 radio sources were detected at $> 5\sigma$. We refer the reader to Smolčić et al. (2017a) for more details.

The VLA 1.4 GHz data are from the VLA-COSMOS Large Project (Schinnerer et al. 2007) and the VLA-COSMOS Deep Project (Schinnerer et al. 2010). The former covered the entire 2 deg^2 COSMOS field with a mean RMS noise of $15 \mu\text{Jy beam}^{-1}$. The latter combined the existing data from the VLA-COSMOS Large Project and an additional VLA L-band observations of the central $50 \text{ arcmin} \times 50 \text{ arcmin}$ subregion in the COSMOS field to reach 1σ sensitivity $\sim 12 \mu\text{Jy beam}^{-1}$ and angular resolution $1''.5 \times 1''.4$. In total, 2856 radio sources with signal-to-noise ratio (SNR) > 5 (peak flux densities) were detected from this combined imaging data. Details are presented in Schinnerer et al. (2010).

2.2.2 GMRT 610 MHz and 325 MHz data

The GMRT 325 MHz data used in this work are from Tisanić et al. (2019). Observations were carried out with GMRT under the project 07SCB01 (PI: C. Crof), which covered the entire $\sim 2 \text{ deg}^2$ COSMOS field and reached a median RMS of $97 \mu\text{Jy beam}^{-1}$ and an angular resolution of $10''.8 \times 9''.5$. In total, 633 radio sources were detected at $> 5\sigma$. Details are presented in Tisanić et al. (2019).

Tisanić et al. (2019) also published GMRT 610 MHz data, which has a median RMS of $39 \mu\text{Jy beam}^{-1}$. However, as described in Tisanić et al. (2019), we find that for sources having measured flux densities at GMRT 325 MHz and 610 MHz, MeerKAT 1.3 GHz, and VLA 1.4 GHz and 3 GHz, there is a deficit in flux density computed at GMRT 610 MHz when compared with flux densities at the other four frequencies. However, it is not a systematic underestimation. Therefore, a simple 20 per cent correction for all sources as employed by some previous work (e.g. Kolokythas et al. 2015) is not appropriate. Tisanić et al. (2019) suggest that this flux density offset might be positionally dependent, but they found that simple exclusion of pointings did not yield a specific pointing which produces these offsets. Since this offset will dramatically affect the measurement of radio spectral index, we exclude the GMRT 610 MHz data in the following analyses.

2.3 Optical/NIR/FIR/submillimeter catalogues

The COSMOS field is one of the largest extragalactic fields with deep multiwavelength data. Previous works have described in detail the available data set in the COSMOS field from X-ray to ultraviolet (UV), optical, near-infrared (NIR), mid-infrared (MIR), FIR, (sub)millimeter and radio bands (e.g. Capak et al. 2007; Sanders et al. 2007; Scoville et al. 2007; Ilbert et al. 2009; Brusa et al. 2010; Civano et al. 2012, 2016; Laigle et al. 2016; Liu et al. 2019; Simpson et al. 2019, 2020). In this work, except the radio data, we do not use these multiwavelength data directly but utilize the published multiwavelength catalogues in the COSMOS field. Here, we give a short introduction of these catalogues.

The optical/NIR photometric catalogue used in this work is the COSMOS2015 catalogue, which is primarily based on the UltraVISTA-DR2 surveys (Laigle et al. 2016). The COSMOS2015

catalogue includes 1182 108 sources detected from a χ^2 sum of $YJHK_s$ and z^{++} imaging data. To match the photometry from Near-UV (NUV) to MIR, Laigle et al. (2016) computed the total fluxes estimated from the corrected 3 arcsec aperture flux, at 31 photometric bands from NUV to MIR, including 12 medium bands. Using the photometrically matched data, Laigle et al. (2016) fitted the SEDs from NUV to MIR for these sources and estimated their photometric redshift, stellar mass, SFR, rest-frame luminosities, and so on. For more details, see Laigle et al. (2016).

The FIR catalogue used in this work is a ‘superdeblended’ FIR to (sub)millimeter photometric catalogue from Jin et al. (2018). They utilized the positions of UltraVISTA K_s - and/or VLA 3 GHz-detected sources (Laigle et al. 2016; Smolčić et al. 2017a) to obtain the point spread function (PSF) fitted flux densities from MIPS $24 \mu\text{m}$ images (Le Floch et al. 2009) and VLA 1.4 GHz and 3 GHz images (Schinnerer et al. 2010; Smolčić et al. 2017a). They then select sources with SNR > 3 at radio or $24 \mu\text{m}$ as priors. With an additional ‘mass-selected’ sample, Jin et al. (2018) adopt a ‘superdeblending’ technique developed by Liu et al. (2018) to ‘deblend’ the FIR to (sub)millimeter photometry from *Spitzer* (Le Floch et al. 2009), *Herschel* (Oliver et al. 2010; Lutz et al. 2011; Béthermin et al. 2012), SCUBA2 (Cowie et al. 2017; Geach et al. 2017), AzTEC (Aretxaga et al. 2011), and MAMBO (Bertoldi et al. 2007) images in the COSMOS field for a total 221 428 sources. Using the ‘superdeblended’ photometry, Jin et al. (2018) fitted FIR to millimeter SEDs for these sources and estimated their SFR by integrating 8–1000 μm infrared luminosities (L_{IR}) derived from the best-fitting SEDs. In this work, we also use the stellar mass from the ‘superdeblended’ catalogue in our analyses (Section 4), which are from Laigle et al. (2016) and Muzzin et al. (2013), and was estimated by using the best-fitting optical to NIR SED with a Chabrier IMF (Chabrier 2003).

The two catalogues of ALMA-detected submillimeter galaxies (SMGs) in the COSMOS field from Liu et al. (2019) and Simpson et al. (2020) are not used in the main analyses of this work but are used in marking this extreme population in our sample of SFGs to compare with previous studies (Thomson et al. 2019; Tisanić et al. 2019). The A³COSMOS³ catalogue includes 1134 ALMA sources with SNR ≥ 5.4 extracted from all publicly available ALMA archive data in the COSMOS field (Liu et al. 2019). The AS2COSMOS catalogue published in Simpson et al. (2020) includes 260 ALMA SMGs detected in observations of an essentially complete sample of 184 SCUBA-2 sources brighter than $S_{850} = 6.2 \text{ mJy}$ (Simpson et al. 2019).

3 ANALYSIS

To investigate the spectral properties of radio-detected sources, we first measure the radio flux densities at different frequencies from VLA, MeerKAT, and GMRT surveys in the COSMOS field. We also describe the selection of SFGs in this section.

3.1 The ‘superdeblended’ flux densities

As described in Section 2.3, Jin et al. (2018) adopted a ‘superdeblending’ technique developed by Liu et al. (2018) to ‘deblend’ FIR to (sub)millimeter photometries of sources selected from NIR or radio. The purpose of the ‘superdeblending’ technique is to deal with source confusion caused by the low angular resolution of FIR

³<https://sites.google.com/view/a3cosmos>

and single-dish (sub)millimeter imaging data when cross-matching them with optical/NIR data. Since the ‘superdeblended’ catalogue we use in this work is from Jin et al. (2018), we take it as an example to shortly describe the basic steps of ‘superdeblending’ technique. First, Jin et al. (2018) used the position of high-resolution K_s -band or VLA 3 GHz-detected sources to perform PSF fitting in MIPS 24 μ m images from Le Floch et al. (2009) and VLA 1.4 and 3 GHz images from Schinnerer et al. (2010) and Smolčić et al. (2017a), respectively, by using the GALFIT (Peng et al. 2002, 2010). For high SNR sources ($\text{SNR} > 10$ at 24 μ m and $\text{SNR} > 20$ at VLA 1.4 GHz and 3 GHz), they run a second pass fitting allowing for up to one and two pixel variations of prior source positions for 24 μ m and radio images, respectively. This returns a cleaner residual image compared to the results of first-pass fitting. Secondly, they performed a Monte Carlo simulation in MIPS 24 μ m, 1.4 GHz, and 3 GHz maps, respectively, to verify potential flux biases and calibrate the uncertainties of photometric measurements. Liu et al. (2018) and Jin et al. (2018) compared the ‘superdeblended’ flux densities with directly measured ones from previous works (e.g. Le Floch et al. 2009; Schinnerer et al. 2010; Smolčić et al. 2017a) and confirmed the reliability of ‘superdeblending’ technique. We refer readers to Liu et al. (2018) and Jin et al. (2018) for more details.

In this work, we aim to match the photometry of the MeerKAT data at 1.3 GHz with those of VLA 1.4 GHz and 3 GHz, and GMRT 325 MHz data. Because of the low resolution of GMRT and MeerKAT data as well as $\sim 10\times$ difference in the angular resolutions between them and VLA 3 GHz data, we also adopt the ‘superdeblending’ technique to measure flux densities at these radio frequencies. To ‘deblend’ the MeerKAT image, we adopt the same $K_s + 3$ GHz prior sample in Jin et al. (2018) and obtain their ‘superdeblended’ flux densities at 1.3 GHz. However, because of the low sensitivity of GMRT data, we only use 3 GHz detection as priors in ‘deblending’ 325 MHz images. This will not affect the sample size of this study as it is based on a radio-selected sample.

To test the self-consistency of the ‘superdeblending’ technique, we compare the ‘superdeblended’ flux densities at 1.3 GHz ($S_{1.3}$) from MIGHTEE survey and at 1.4 GHz ($S_{1.4}$) from VLA-COSMOS survey for the 1599 sources with $\text{SNR} > 3$ at both frequencies as shown in Fig. 1. We limit our analyses within the area of the MeerKAT primary beam (1 deg diameter) to avoid low sensitivity regions at the edges of MIGHTEE survey. The relatively lower median value of ($S_{1.3}/S_{1.4}$) within the lowest 1.4 GHz flux-density bin (Fig. 1) is caused by the very different sensitivities of the two data sets (Section 2). For the cut of $\text{SNR} > 3$, the low sensitivity of VLA 1.4 GHz data results in a relatively larger uncertainty of measured flux density at 1.4 GHz (Fig. 1), which, in return, leads to a relatively higher flux density cut at VLA 1.4 GHz. In contrast, the higher SNR cut ($\text{SNR} > 5$) is less affected by this effect, as shown in Fig. 1. We therefore mainly rely on the sample with $\text{SNR} > 5$ at both frequencies in verifying the self-consistency of the ‘superdeblending’ technique.

The inset plot in Fig. 1 shows the distributions of normalized difference between the two flux densities for the sources with $\text{SNR} > 3$ and > 5 , respectively. The normalized difference is defined as

$$\delta = (S_{1.4} - S_{1.3}) / \sqrt{(\sigma_{1.4}^2 + \sigma_{1.3}^2)}, \quad (1)$$

where σ refers to the uncertainty of ‘superdeblended’ flux density. The mean differences are $\delta = 0.02$ and -0.18 with the scatters of $\sigma_\delta = 0.94$ and 0.96 for the two subsamples, respectively. The negative δ and the median ($S_{1.3}/S_{1.4}$) = $1.04^{+0.02}_{-0.01}$ for sources with $\text{SNR} > 5$ might be indicative of the slightly different effective frequencies of the two data sets, which results in ($S_{1.3}/S_{1.4}$) = 1.06 if we assume

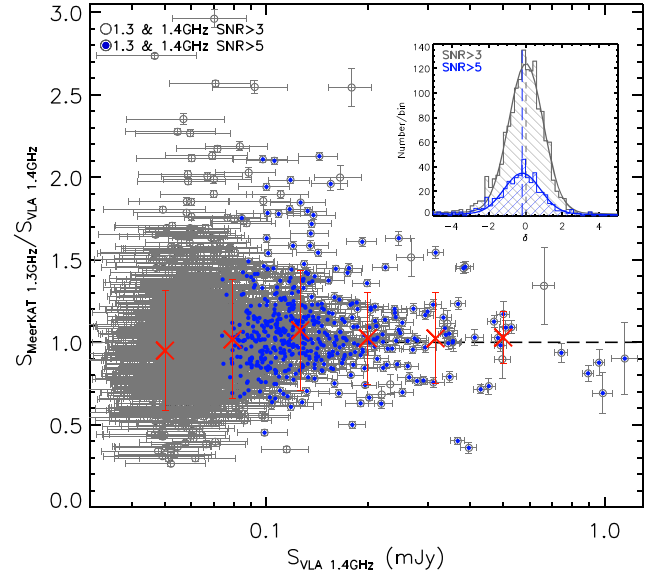


Figure 1. Ratio of ‘superdeblended’ flux densities at MeerKAT 1.3 GHz and VLA 1.4 GHz as a function of 1.4 GHz flux density for the 1599 sources with $\text{SNR} > 3$ at both frequencies. We mark the sources that have $\text{SNR} > 5$ at both frequencies by blue solid points. The red crosses are the median ($S_{1.3}/S_{1.4}$) for sources within each 1.4 GHz flux density bin. The inset plot shows the distribution of normalized difference of ‘superdeblended’ flux densities at these two frequencies for sources with $\text{SNR} > 3$ and $\text{SNR} > 5$ at both frequencies, respectively. Overall, despite the scatter, the ‘superdeblended’ flux densities extracted from the VLA 1.4 GHz and MeerKAT 1.3 GHz data are consistent with each other within the uncertainties for sources with $\text{SNR} > 5$ if we consider their slightly different effective frequencies.

the spectral index $\alpha = -0.8$ between 1.3 and 1.4 GHz. Therefore, although some of the sources appear to be outliers, the flux densities at these two frequencies are consistent with each other for the sources with $\text{SNR} > 5$.

3.2 SFGs selection

Due to the high resolution and the high sensitivity of VLA 3 GHz data, we select sources with $\text{SNR} > 5$ at 3 GHz from the ‘superdeblended’ catalogue (Jin et al. 2018) to study the radio spectral properties in this work. The 3 GHz is also the highest observer-frame frequency we use in studying radio spectrum in this work. Compared with selecting sources at low frequency, which bias the sample to the sources with a relatively steeper radio spectrum, the choice of 3 GHz as the selection frequency provides a less-biased sample for studying the radio spectrum of SFGs in this work. In addition, the high sensitivity of MeerKAT 1.3 GHz data guarantees a high completeness of studying radio spectrum at 1.3–3 GHz for 3 GHz-selected SFGs. We will come back to this in Section 4.

Within the area of the MeerKAT primary beam (1 deg diameter), there are 4509 VLA 3 GHz sources with $\text{SNR} > 5$. We first select the star formation dominated galaxies by removing previously classified AGNs or red galaxies from these radio sources.

For X-ray AGNs in the COSMOS field, Marchesi et al. (2016) cross-correlated the *Chandra* X-ray detected sources from the *Chandra* COSMOS legacy Survey (Civano et al. 2016) with optical/NIR sources in the literature and identified optical/NIR counterparts for 97 per cent of 4016 *Chandra* X-ray sources. We first cross-match VLA 3 GHz radio sources with optical/NIR counterparts of X-ray sources from Marchesi et al. (2016) by using a matching radius of

1 arcsec. We find that 566/4509 radio sources have *Chandra* X-ray detections and 519 of them are classified as AGNs in Marchesi et al. (2016) according to their rest-frame X-ray luminosity ($L_X > 10^{42} \text{ erg s}^{-1}$). We also cross-correlate these 3 GHz radio sources with optical/NIR counterparts of 2012 X-ray sources from the *XMM-Newton* X-ray survey (Cappelluti et al. 2009; Brusa et al. 2010) and find an additional six X-ray detected AGNs.

The X-ray AGN-selection tends to miss low-luminosity AGNs at high redshift, as they typically do not have accretion-related X-ray emission (Marchesi et al. 2016). Therefore, we also adopt the MIR colour–colour selection criteria in Donley et al. (2012) for $z \leq 3$ sources. Since the ‘superdeblended’ catalogue includes all VLA 3 GHz sources with $\text{SNR} > 5$ (Jin et al. 2018), we use the matched photometry from this catalogue and find that 2760 sources in our sample have $\text{SNR} > 3$ in all four IRAC bands. Among these, 240 meet the selection criteria of MIR AGNs (Donley et al. 2012).

In the ‘superdeblended’ catalogue, Jin et al. (2018) classified 726/4509 radio-excess sources as radio-loud AGNs. The criteria for radio excess are defined as $(S_{\text{OBS,radio}} - S_{\text{SED,radio}}) / \sqrt{(\sigma_{\text{OBS,radio}}^2 + \sigma_{\text{SED,radio}}^2)} > 3$ and $S_{\text{OBS,radio}} > 2 \times S_{\text{SED,radio}}$ (Liu et al. 2018), where $S_{\text{OBS,radio}}$ is the observed radio flux density and $S_{\text{SED,radio}}$ is the flux density predicted by the best-fitting FIR to millimeter SEDs (Section 2.3). For the 566 and 311 VLA 3 GHz radio sources that are cross-matched with the optical/NIR counterparts of X-ray sources from Marchesi et al. (2016) and from Brusa et al. (2010), 160 of them are identified as AGNs according to their best-fitting optical to NIR SED in Marchesi et al. (2016), and 392 and 183 are spectroscopically identified AGNs in Marchesi et al. (2016) and Brusa et al. (2010), respectively. In addition, 3958/4509 radio sources are also included in the 5σ 3 GHz catalogue presented in Smolčić et al. (2017b). Among them, 501, 271, 578, and 793 are classified as X-ray, MIR, SED, and radio excess AGNs in Smolčić et al. (2017b). By integrating these classifications from the literature, in total we remove 1916 AGNs from the 4509 VLA 3 GHz radio sources.

To further ensure a clean sample of SFGs, we also remove red galaxies with both $(M_{\text{NUV}} - M_r) > 3.5$ and non-detection in the *Herschel* bands (100, 160, 250, 350, and 500 μm) from our sample, since previous work proved that this population shows typical properties that are consistent with those of radio AGN host galaxies (e.g. Best et al. 2006; Sadler et al. 2014; Smolčić et al. 2017b). The absolute rest-frame *GALEX* NUV and Subaru *r*-band magnitudes are taken from the COSMOS2015 catalogue (Laigle et al. 2016), while the non-detection by *Herschel* is defined as sources without $> 3\sigma$ detection at any *Herschel* band. There are 761 galaxies that meet these selection criteria and 573 of them overlap with our AGN sample. Therefore, after removing 2104 AGN and/or red galaxies, we obtain a sample of 2405 SFGs with $\text{SNR} > 5$ at VLA 3 GHz. Among these, 2094 have a photometric redshift or spectroscopic redshift in the ‘superdeblended’ catalogue, in which the photometric redshift is from the COSMOS2015 catalogue (Laigle et al. 2016) while the spectroscopic redshift is from a new COSMOS master spectroscopic catalogue (Jin et al. 2018; M. Salvato et al., in preparation).

4 RESULTS

To study the radio spectral properties of SFGs, we first extract their flux densities from MeerKAT 1.3 GHz, and GMRT 325 MHz maps at their 3 GHz positions by applying the ‘superdeblending’ technique (Jin et al. 2018).

4.1 Spectral index of SFGs between the observer-frame frequencies of 1.3–3 GHz

For the 2,094 VLA 3 GHz-selected SFGs, 96 per cent (2012/2094) of them have measured ‘superdeblended’ flux densities with $\text{SNR} > 1$ at MeerKAT 1.3 GHz. Among them, 1851 (92 per cent) have $\text{SNR} > 3$ while the rest 161 have $1 < \text{SNR} \leq 3$. However, because of the low sensitivity of VLA 1.4 GHz data, although 85 per cent (1783/2094) of these 3 GHz-selected SFGs have the measured flux densities at VLA 1.4 GHz, only a further 42 per cent (756/1783) have $\text{SNR} > 3$. Therefore, in this work, we use the ‘superdeblended’ flux densities at MeerKAT 1.3 GHz and VLA 3 GHz to measure the radio spectral index between the observer-frame frequencies of 1.3–3 GHz for our selected SFGs. The radio spectral index, $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$, is measured by performing a linear fit in logarithmic space, and its uncertainty is inherited from the uncertainties of the flux densities. The flux densities at VLA 1.4 GHz are included in the fit if the SFGs also have $\text{SNR} > 3$ at VLA 1.4 GHz.

There are 712 VLA 3 GHz-selected SFGs with $\text{SNR} > 3$ at both MeerKAT 1.3 GHz and VLA 1.4 GHz. We measure their radio spectral index between the observer-frame frequencies of 1.3–3 GHz and obtain a median radio spectral index of $\alpha_{1.3\text{GHz}}^{3\text{GHz}} = -0.80 \pm 0.01$ and $\alpha_{1.3\text{GHz}}^{3\text{GHz}} = -0.79 \pm 0.01$ if we include or exclude the VLA 1.4 GHz data in the measurement, respectively. Here, the error of the median spectral index is estimated from bootstrap resampling. The scatters of both measurements are $\sigma = 0.31$. These results further suggest the self-consistency of ‘superdeblending’ technique (Section 3.1).

For the 2012 SFGs that have measured flux densities at MeerKAT 1.3 GHz ($\text{SNR} > 1$), their median spectral index is $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} = -0.78 \pm 0.01$ with a scatter of $\sigma = 0.37$. However, if we limit our sample to the SFGs with $\text{SNR} > 3$ at MeerKAT 1.3 GHz, we obtain a median spectral index of $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} = -0.80 \pm 0.01$ and a scatter of $\sigma = 0.35$. Therefore, a higher SNR cut at MeerKAT 1.3 GHz will bias our sample towards SFGs with a relatively steeper radio spectrum at observer-frame 1.3–3 GHz, but we find that this effect is negligible because of the high sensitivity of MeerKAT 1.3 GHz data. We thus study the correlations between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and physical properties of SFGs based on the sample of SFGs with $\text{SNR} > 3$ at MeerKAT 1.3 GHz as shown in Fig. 2. For the 161 SFGs with $1 < \text{SNR} \leq 3$ at MeerKAT 1.3 GHz, because of the relatively larger uncertainties of their 1.3 GHz flux densities, we use the 5σ limit of MeerKAT data to estimate the lower limit of their radio spectral indices (Fig. 2).

In Fig. 2, we show the measured spectral index between the observer-frame frequencies of 1.3–3 GHz as a function of corresponding rest-frame frequency of VLA 3 GHz, redshift, stellar mass, SFR, and specific star formation rate (sSFR) for the 2012 SFGs. The sSFR is defined as the ratio of SFR and stellar mass. The relatively steeper spectrum at the lowest redshift, sSFR, and the two lowest SFR bins is likely to be caused by missing flux in the high-resolution VLA 3 GHz observations. This results from the fact that for extended galaxies at low redshift, only part of their flux densities have been recovered in the high-resolution 3 GHz observations. Therefore, we simply exclude the SFGs with $z < 0.5$ in the linear fits shown in Fig. 2.

As shown in the first panel (a1) of Fig. 2, on average, the spectral index at observer-frame 1.3–3 GHz is not strongly correlated with rest-frame frequency (therefore also redshift) of SFGs. The slope of the linear fit between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and rest-frame frequency ($\nu_{\text{rest}, 3\text{GHz}}$) is $\beta_{(\alpha, \nu_{\text{rest}, 3\text{GHz}})} = 0.002 \pm 0.002$. We also show the trends between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and rest-frame frequency for SFGs with 10.3

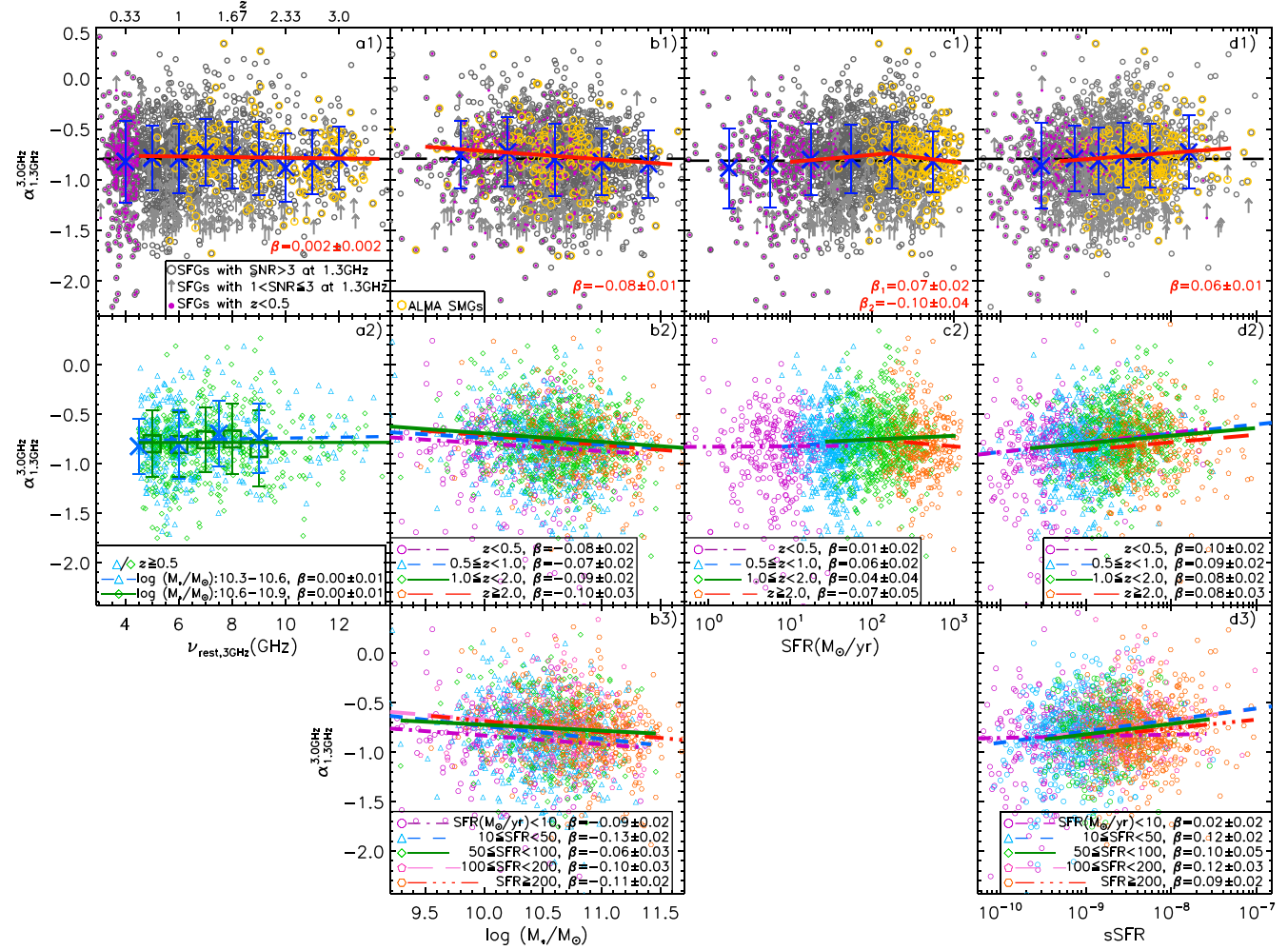


Figure 2. Spectral index between the observer-frame frequencies of 1.3–3 GHz, $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$, as a function of rest-frame frequency corresponding to observed 3 GHz, redshift (a1), stellar mass (b1), SFR (c1), and sSFR (d1) for the 2012 VLA 3 GHz-selected SFGs. The upward arrows show the lower limit of $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ for the 161 SFGs with $1 < \text{SNR} \leq 3$ at MeerKAT 1.3 GHz. The black dashed line represents the median radio spectral index of $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} = -0.80 \pm 0.01$ for the remaining 1851 SFGs with $\text{SNR} > 3$ at MeerKAT 1.3 GHz. The blue cross or green square represents the median $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ of SFGs in each bin while the error bar shows the standard deviation. The slightly steeper spectrum at the lowest redshift (rest-frame), sSFR, and the two lowest SFR bins might be biased by the missing flux of high-resolution VLA 3 GHz observations for extended sources at low redshift. We therefore exclude galaxies with $z < 0.5$ in the linear fits of the correlation between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and a physical property of the galaxies within stellar mass (a2), redshift (b2, c2, and d2) or SFR (b3, and d3) bins. The plot c2 suggests that the trend between SFR and $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ is caused by the redshift-dependent selection effect of our radio flux-limited sample. Therefore, the trend between radio spectral index and sSFR (d1) might be another manifestation of the correlation between the radio spectrum and the stellar mass of SFGs. The plots in the first and second columns suggest that the trends between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and redshift or stellar mass are very unlikely to be caused by selection effect. Overall, on average, the radio spectral slope at observer-frame 1.3–3 GHz is not strongly correlated with redshift/rest-frame frequency, but slightly steepens with increasing stellar mass and flattens with sSFR for radio-selected SFGs, although the scatter is large.

$\leq \log(M_*/M_\odot) < 10.6$ and $10.6 \leq \log(M_*/M_\odot) < 10.9$, respectively (plot a2 of Fig. 2) and confirm that this trend is independent from the stellar mass of SFGs. Galaxies with $z < 0.5$ are excluded in the plot a2 of Fig. 2 to avoid the effect from missing flux of VLA 3 GHz observations.

The second panel (b1) of Fig. 2 shows that the radio spectral slope at observer-frame 1.3–3 GHz steepens slightly with increasing stellar mass with a linear fitted slope of $\beta = -0.08 \pm 0.01$. We find that this trend can be observed within different redshift or SFR bins as shown in the plot b2 and b3 of Fig. 2.

The third panel (c1) of Fig. 2 shows that, on average, the radio spectral slope first slightly flattens with increasing SFR but reverses from $\text{SFR} \gtrsim 200 M_\odot \text{yr}^{-1}$. The slopes of the linear fits are $\beta_{(\alpha, \text{SFR} \lesssim 200 M_\odot \text{yr}^{-1})} = 0.07 \pm 0.02$ and $\beta_{(\alpha, \text{SFR} \gtrsim 200 M_\odot \text{yr}^{-1})} = -0.10 \pm 0.04$, respectively. However, the plot c2 of Fig. 2 suggests that this is caused by redshift-dependent selection effect of our radio flux-limited sample. These analyses suggest that the trend between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and sSFR ($\beta_{(\alpha, \text{sSFR})} = 0.06 \pm 0.01$, plot d1 of Fig. 2) is to a large extent another manifestation of the correlation between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and stellar mass (b1 of Fig. 2).

We also include SFGs with $z < 0.5$ in the linear fits of the correlations between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and the physical properties shown in Fig. 2. The slopes of these linear fits change to $\beta_{(\alpha, v_{\text{rest}}, 3\text{GHz})} = 0.005 \pm 0.002$, $\beta_{(\alpha, \log(M_*/M_\odot))} = -0.06 \pm 0.01$, $\beta_{(\alpha, \text{sSFR} \lesssim 200 M_\odot \text{ yr}^{-1})} = 0.05 \pm 0.01$, and $\beta_{(\alpha, \text{sSFR})} = 0.07 \pm 0.01$, respectively. These results suggest that the missing flux of high-resolution VLA 3 GHz observations for extended sources at low redshift slightly affects the study of the correlation between $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}}$ and a physical property of the SFGs in this work.

We mark the 187 SFGs with ALMA detection from AS2COSMOS (Simpson et al. 2020) or A3COSMOS surveys (Liu et al. 2019) in Fig. 2 and find that their radio spectra at observer-frame 1.3–3 GHz show the statistical properties consistent with the entire sample.

Overall, out of the five physical properties shown in Fig. 2, only the stellar mass and sSFR show robust correlations with the radio spectral index at observer-frame 1.3–3 GHz of SFGs, although the scatter is large. We will discuss the physical mechanisms responsible for the trends between radio spectral index and these physical properties in Section 5.1.

4.2 Spectral indices of SFGs between the observer-frame frequencies of 0.33–3 GHz

To study the radio spectral properties at low frequency, we include the GMRT 325 MHz data to measure the radio spectral index between the observer-frame frequencies of 0.33–3 GHz for selected SFGs in the COSMOS field.

4.2.1 Comparison of radio spectrum at high and low frequencies

For the 2094 VLA 3 GHz-selected SFGs, we have measured the spectral index at observer-frame 1.3–3 GHz for 96 per cent (2012/2094) of them (Section 4.1). However, because of the low sensitivity of the GMRT 325 MHz data, only half of these 3 GHz-selected SFGs have a measured ‘superdeblended’ flux density with $\text{SNR} > 1$ at GMRT 325 MHz. Therefore, we apply an additional flux density cut of $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ to obtain a less-biased sample of SFGs for studying the radio spectral properties at observer-frame 0.33–3 GHz. The criterion $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ is chosen in order to balance the completeness of the SFGs with measured flux density at GMRT 325 MHz and the sample size. Within the region for our analyses, there are 220 SFGs with $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$. Among them, 91 per cent (201/220) and 84 per cent (184/220) have $\text{SNR} > 1$ at MeerKAT 1.3 GHz and GMRT 325 MHz, respectively. For the 19 sources that are brighter than $50 \mu\text{Jy}$ at VLA 3 GHz but lack a measured ‘superdeblended’ flux density at MeerKAT 1.3 GHz, we visually inspect the MeerKAT 1.3 GHz image and the residual map. We find that their residual maps are not clean after two rounds of GALFIT PSF fitting (Jin et al. 2018), which indicates that their flux density profiles cannot be well-approximated by the Gaussian distribution. For this reason, their MeerKAT 1.3 GHz flux densities cannot be accurately derived by the ‘superdeblending’ technique. Likewise, of the 36 sources that lack the measured flux density with $\text{SNR} > 1$ at GMRT 325 MHz, four of them are also caused by the inability to fit with a Gaussian profile. In addition, the other 32 simply have $\text{SNR} < 1$ at GMRT 325 MHz.

We first measure the radio spectral indices between the observer-frame frequencies of 0.33–1.3 GHz, 0.33–3 GHz, and 1.3–3 GHz for a sample of 166, 184, and 201 SFGs, respectively, using the two flux densities at the endpoints of each frequency range. The median radio spectral indices are $\alpha_{325\text{MHz}}^{1.3\text{GHz}} = -0.59^{+0.02}_{-0.03}$, $\alpha_{325\text{MHz}}^{3\text{GHz}} = -0.65^{+0.02}_{-0.03}$, and $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} = -0.73^{+0.01}_{-0.02}$ with a scatter of $\sigma = 0.29, 0.19$, and

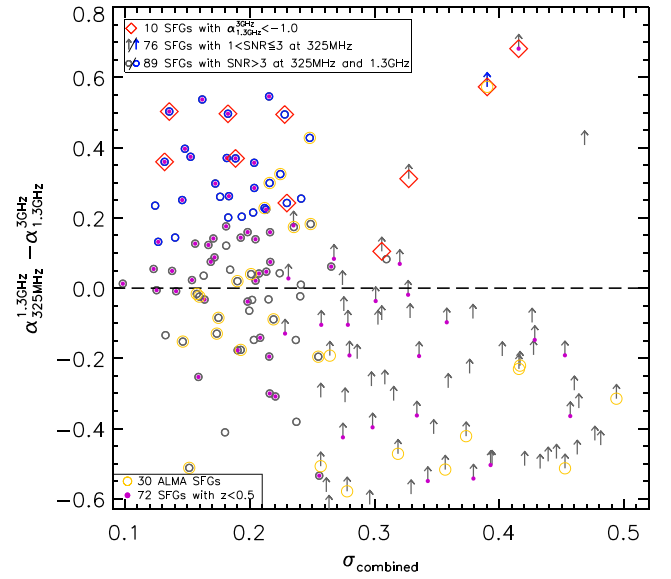


Figure 3. Difference between radio spectral indices at low and high frequencies ($\alpha_{325\text{MHz}}^{1.3\text{GHz}} - \alpha_{325\text{MHz}}^{3\text{GHz}}$) as a function of their combined uncertainties for the 166 SFGs with $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ and measured flux densities at the three radio frequencies. The upward arrows show the lower limit of ($\alpha_{325\text{MHz}}^{1.3\text{GHz}} - \alpha_{325\text{MHz}}^{3\text{GHz}}$) for the 76 SFGs with $1 < \text{SNR} \leq 3$ at GMRT 325 MHz. The blue symbols are the 30 SFGs, which have a low-frequency radio spectrum significantly flatter than that at high frequency. We also mark the 72 SFGs with $z < 0.5$ (purple points), the 30 ALMA SMGs (yellow circles), and 10 SFGs with $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} < -1.0$ (red rhombs), respectively.

0.22, respectively. Therefore, on average, the radio spectrum of SFGs flattens at low frequency.

For the SFGs with $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ and measured flux densities at MeerKAT 1.3 GHz, 99 per cent (198/201) of them have $\text{SNR} > 3$ at MeerKAT 1.3 GHz. However, for the 184 SFGs with measured flux density at GMRT 325 MHz, only 58 per cent (107/184) of them have $\text{SNR} > 3$ at GMRT 325 MHz. The measured radio spectral indices for a sample of 89, 107, and 198 SFGs with $\text{SNR} > 3$ at the two corresponding frequencies are $\alpha_{325\text{MHz}}^{1.3\text{GHz}} = -0.69^{+0.03}_{-0.06}$, $\alpha_{325\text{MHz}}^{3\text{GHz}} = -0.74^{+0.01}_{-0.03}$, and $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} = -0.73^{+0.01}_{-0.02}$ with a scatter of $\sigma = 0.19, 0.13$, and 0.22 , respectively. Therefore, a higher SNR cut at GMRT 325 MHz will bias our sample towards the SFGs with a relatively steeper radio spectrum at observer-frame 0.33–3 GHz, as one would expect. These suggest that high-sensitivity multifrequency data are necessary for a complete study of radio spectrum of SFGs.

We also compare the measured radio spectral index at high and low frequencies for the 166 SFGs with measured flux densities at the three radio frequencies individually (Fig. 3). For the 76 SFGs that have $1 < \text{SNR} \leq 3$ at GMRT 325 MHz, we use the 5σ limit of GMRT 325 MHz data to estimate the lower limit of their $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$. In addition, for the 90 SFGs that have $\text{SNR} > 3$ at GMRT 325 MHz, one of them has $\text{SNR} = 2$ at MeerKAT 1.3 GHz. Therefore, in total, there are 89 SFGs with $\text{SNR} > 3$ at both MeerKAT 1.3 GHz and GMRT 325 MHz. We find that 33 per cent (29/89) of them have the difference between radio spectral indices at low and high frequencies ($\alpha_{325\text{MHz}}^{1.3\text{GHz}} - \alpha_{325\text{MHz}}^{3\text{GHz}}$) larger than the combined uncertainty, which is estimated as

$$\sigma_{\text{combined}} = \sqrt{\left(\sigma_{\alpha_{325\text{MHz}}^{1.3\text{GHz}}}^2 + \sigma_{\alpha_{325\text{MHz}}^{3\text{GHz}}}^2\right)}. \quad (2)$$

In addition, one SFG has the estimated lower limit of ($\alpha_{325\text{MHz}}^{1.3\text{GHz}} - \alpha_{325\text{MHz}}^{3\text{GHz}}$) larger than its σ_{combined} . We also mark the 10 SFGs that have $\alpha_{1.3\text{GHz}}^{3.0\text{GHz}} < -1$ in Fig. 3. We find that five out of these 10 have $z <$

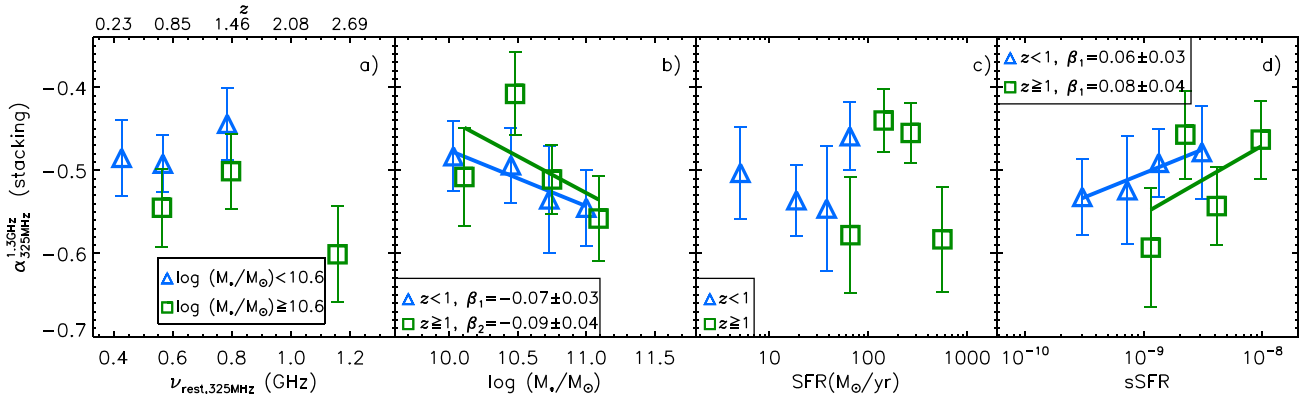


Figure 4. Averaged $\alpha_{1.3\text{GHz}/325\text{MHz}}^{1.3\text{GHz}}$, which is estimated by stacking MeerKAT and GMRT images for all of the 2,094 VLA 3 GHz-selected SFGs, as functions of rest-frame frequency corresponding to observed GMRT 325 MHz, redshift, stellar mass, SFR, and sSFR. We divide the sample into the massive ($\log(M_*/M_\odot) \geq 10.6$) and less-massive ($\log(M_*/M_\odot) < 10.6$) subsamples when stacking SFGs within different redshift bins and find that, on average, the massive SFGs have a relatively steeper radio spectrum. We also divide our sample into high-redshift ($z \geq 1$) and low-redshift ($z < 1$) subsamples when stacking SFGs within different stellar mass, SFR, and sSFR bins. The results show that the averaged radio spectral slope at low frequency also steepens with increasing stellar mass, and flattens with increasing sSFR despite some fluctuations at high redshift ($z \geq 1$).

0.5, while four have $z = 0.7\text{--}0.8$. Therefore, some of these sources' 3 GHz flux densities are very likely affected by the missing flux of VLA observations. For SFGs brighter than $50\ \mu\text{Jy}$ at 3 GHz, we thus estimate that > 15 per cent of them have a significantly flatter radio spectrum at low frequency.

4.2.2 Stacking all of the 3 GHz-selected SFGs

For all of the 2094 VLA 3 GHz selected SFGs in our sample, only 195 of them have $\text{SNR} > 3$ at GMRT 325 MHz. In addition, the cut of $\text{SNR} > 3$ at GMRT 325 MHz biases our sample towards the SFGs with relatively steeper radio spectrum observed at 0.33–3 GHz as described in Section 4.2.1. We therefore perform a stacking analysis to study the radio spectral properties at observer-frame 0.33–1.3 GHz.

For these 2094 SFGs, we stack both GMRT 325 MHz and MeerKAT 1.3 GHz imaging data at their 3 GHz positions to measure the averaged radio spectral index at observer-frame 0.33–1.3 GHz. We use the peak flux densities of the average-stacked images to estimate the radio spectral index at low frequency and show the results in Fig. 4 and Table 1. To compare the correlations of radio spectral index and physical properties of galaxies shown in Fig. 2, we stack the 3 GHz-selected SFGs by redshift, stellar mass, SFR, and sSFR bins, respectively.

As shown in Fig. 4 and Table 1, when stacking the SFGs within different redshift bins, we further divide the sample into massive ($\log(M_*/M_\odot) \geq 10.6$) and less-massive ($\log(M_*/M_\odot) < 10.6$) subsamples. We find that the averaged radio spectral slope at low-frequency slightly steepens at $\nu_{\text{rest},325\text{MHz}} \gtrsim 1.0\text{ GHz}$ (plot a of Fig. 4). However, as shown in Fig. 4, our sample at $z > 2$ are dominated by massive SFGs. Therefore, the results suggest that, on average, massive SFGs have a relatively steeper radio spectrum at observer-frame 0.33–1.3 GHz.

For the 2094 SFGs, 2090 of them have estimated stellar mass from the ‘superdeblended’ catalogue (Jin et al. 2018). We further divide them into high-redshift ($z \geq 1$) and low-redshift ($z < 1$) subsets to investigate the selection effect of our radio flux-limited sample. Except the lowest stellar mass bin, which is affected the most by the incompleteness of our flux-limited sample, the results from stacking show that the averaged radio spectral slope at low frequency also steepens with increasing stellar mass (Fig. 4).

We also divide our sample into high-redshift ($z \geq 1$) and low-redshift ($z < 1$) subsets when stacking SFGs within different SFR bins and confirm that the trend between the averaged radio spectral index and SFR is caused by the redshift-dependent selection effect of our radio flux-limited sample (plot c of Fig. 4). As shown in the plot (d) of Fig. 4, the averaged radio spectral slope at low frequency also flattens with increasing sSFR in general, although the high-redshift subsample exhibits the greater fluctuation in the trend.

For the 2012 SFGs with measured radio spectral index at high frequency ($\alpha_{1.3\text{GHz}}^{3\text{GHz}}$), we stack them by high-frequency radio spectral index bins and show the results in Table 1. We find that the averaged radio spectral index at low frequency is positively correlated with that at high frequency, i.e. the SFGs within the lowest $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ bin also have the steepest averaged radio spectrum at low frequency, and vice versa. We further divide the 2012 SFGs into high-redshift ($z \geq 1$) and low-redshift ($z < 1$) subsamples and find that the positive correlation between the averaged radio spectral indices at low and high frequencies is present in both redshift subsamples despite some fluctuations.

Overall, the stacking analyses of all the 3 GHz-selected SFGs suggest that the trends between radio spectral index and physical properties of SFGs are consistent at low and high frequencies. In addition, the averaged radio spectral index at low and high frequencies show a positive correlation.

4.3 FIR-radio correlation

With the measured radio spectral indices of SFGs, we study the correlation between the infrared and radio emission of SFGs and investigate how different assumptions of k -correction affect this. In this work, we use rest-frame radio luminosities at 1.3 GHz ($L_{1.3\text{GHz}}$) in the FIRRC of SFGs to be consistent with the observational frequency of our MIGHTEE survey.

4.3.1 FIR-radio correlation for bright 3 GHz SFGs

To measure $L_{1.3\text{GHz}}$, the spectral index at low frequency ($\alpha_{1.3\text{GHz}}^{1.3\text{GHz}}$) is essential because the emission at rest-frame 1.3 GHz from a galaxy with $z > 0$ is shifted into $\nu_{\text{obs}} < 1.3\text{ GHz}$. Therefore, the appropriate spectral index to use in k -correction should be $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$. However,

Table 1. Results from stacking the 3 GHz-selected SFGs.

Redshift (z)	z < 0.5	0.5–1.0	1.0–2.0	≥ 2.0
Number	373	668	703	350
S ₃₂₅ MHz *	203 ± 9	125 ± 5	109 ± 4	109 ± 7
S _{1.3} GHz *	99 ± 4	62 ± 2	56 ± 1	50 ± 1
α ₃₂₅ MHz	−0.51 ± 0.04	−0.51 ± 0.03	−0.48 ± 0.03	−0.56 ± 0.05
log(M _* /M _⊙) < 10.6			log(M _* /M _⊙) ≥ 10.6	
z < 0.5	0.5–1.0	1.0–2.0	z < 1.0	1.0–2.0
279	395	284	364	418
172 ± 9	106 ± 4	84 ± 4	191 ± 10	126 ± 7
88 ± 4	54 ± 2	46 ± 1	90 ± 3	63 ± 2
−0.49 ± 0.05	−0.49 ± 0.04	−0.44 ± 0.04	−0.55 ± 0.04	−0.50 ± 0.05
log(M _* /M _⊙)	< 10.3	10.3–10.6	10.6–10.9	≥ 10.9
Number	493	550	578	469
S ₃₂₅ MHz *	112 ± 5	116 ± 4	141 ± 6	152 ± 8
S _{1.3} GHz *	57 ± 2	62 ± 2	69 ± 2	70 ± 2
α ₃₂₅ MHz	−0.49 ± 0.03	−0.46 ± 0.03	−0.52 ± 0.04	−0.56 ± 0.04
z < 1			z ≥ 1	
log(M _* /M _⊙) < 10.3	10.3–10.6	10.6–10.9	≥ 10.9	log(M _* /M _⊙) < 10.3
353	321	242	122	140
125 ± 6	142 ± 6	188 ± 13	196 ± 14	78 ± 6
64 ± 2	72 ± 3	90 ± 5	90 ± 4	39 ± 2
−0.48 ± 0.04	−0.49 ± 0.05	−0.54 ± 0.06	−0.55 ± 0.05	−0.51 ± 0.06
SFR (M _⊙ yr ^{−1})	z < 1			
SFR < 10	SFR < 10	10–30	30–50	≥ 50
239	239	348	251	203
159 ± 9	159 ± 9	146 ± 7	138 ± 12	176 ± 8
79 ± 4	79 ± 4	69 ± 3	65 ± 4	93 ± 4
−0.50 ± 0.05	−0.50 ± 0.05	−0.54 ± 0.04	−0.55 ± 0.08	−0.46 ± 0.04
log(M _* /M _⊙) < 10.3	10.3–10.6	10.6–10.9	≥ 10.9	log(M _* /M _⊙) < 10.3
336	347	111 ± 6	137 ± 8	137 ± 8
55 ± 2	63 ± 2	−0.51 ± 0.04	−0.56 ± 0.05	−0.51 ± 0.04
100–200	200–400	≥ 400		
301	297	201		
90 ± 4	110 ± 5	167 ± 13		
49 ± 1	59 ± 2	74 ± 3		
−0.44 ± 0.04	−0.46 ± 0.04	−0.58 ± 0.07		−0.58 ± 0.06

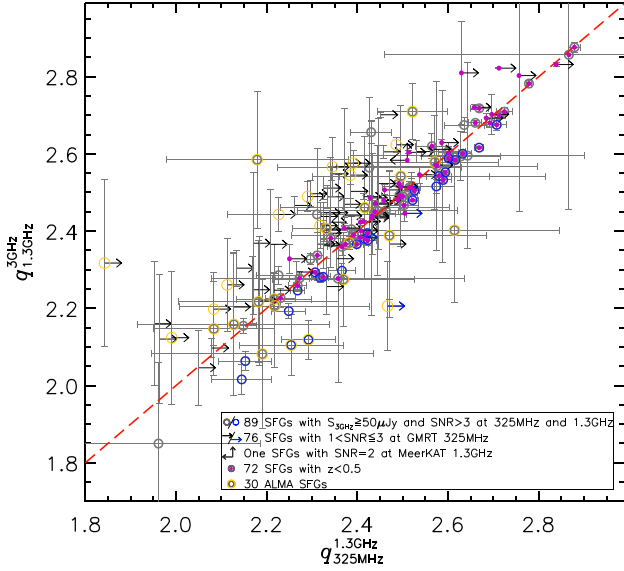


Figure 5. Comparison of q_{IR} based on radio spectral indices at low and high frequencies for the 166 SFGs with $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ and measured flux densities at the three radio frequencies. The $q_{325\text{MHz}}^{1.3\text{GHz}}$ is measured by adopting $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ to k -correct the observed flux density, while $q_{1.3\text{GHz}}^{3\text{GHz}}$ is based on $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ -derived k -correction. The circles are the 89 SFGs with $\text{SNR} > 3$ at both GMRT 325 MHz and MeerKAT 1.3 GHz. For the 76 SFGs with $1 < \text{SNR} \leq 3$ at GMRT 325 MHz, we use the lower limit of their radio spectral indices to estimate the lower limit of their $q_{325\text{MHz}}^{1.3\text{GHz}}$ as shown by rightward arrows. The remaining one SFG has $\text{SNR} > 3$ at GMRT 325 MHz but has $\text{SNR} = 2$ at MeerKAT 1.3 GHz. We show its upper limit of $q_{325\text{MHz}}^{1.3\text{GHz}}$ and lower limit of $q_{1.3\text{GHz}}^{3\text{GHz}}$ by leftward and upward arrows. We also mark the SFGs with the measured $(q_{325\text{MHz}}^{1.3\text{GHz}} - q_{1.3\text{GHz}}^{3\text{GHz}})$ larger than their combined uncertainties by blue symbols. Errors of q_{IR} are derived from the uncertainties of infrared and radio luminosities.

as we described in Section 4.2, because of the low sensitivity of the GMRT data, only the sample of the SFGs brighter than $50 \mu\text{Jy}$ at 3 GHz has a relatively complete (> 80 per cent) measurement of flux densities at GMRT 325 MHz. We thus first use $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ to k -correct the observed flux density at MeerKAT 1.3 GHz and derive the rest-frame luminosity at 1.3 GHz for the 89 SFGs with $\text{SNR} > 3$ at both GMRT 325 MHz and MeerKAT 1.3 GHz. For the 76 SFGs with $1 < \text{SNR} \leq 3$ at GMRT 325 MHz, we use the lower limit of their $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$. The remaining one SFG has $\text{SNR} > 3$ at GMRT 325 MHz but has $\text{SNR} = 2$ at MeerKAT 1.3 GHz. We therefore use the lower limit of its $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ to estimate the upper limit of $q_{325\text{MHz}}^{1.3\text{GHz}}$ (Fig. 5). Combining with their total infrared luminosities from the ‘superdeblended’ catalogue (Jin et al. 2018), we measure the logarithmic infrared-to-radio luminosity ratio (e.g. Helou et al. 1985; Condon 1992; Ivison et al. 2010):

$$q_{\text{IR}} = \log \left[\frac{L_{\text{IR}}}{(3.75 \times 10^{12} \text{ W})} \times \frac{\text{WHz}^{-1}}{L_{1.3\text{GHz}}} \right]. \quad (3)$$

We obtain a median $q_{\text{IR}} = 2.46_{-0.04}^{+0.03}$ and a scatter of $\sigma = 0.18$ for the 89 SFGs with $\text{SNR} > 3$ at both GMRT 325 MHz and MeerKAT 1.3 GHz. To investigate how different k -corrections affect the FIRRC of SFGs, we also use the spectral indices at high frequency ($\alpha_{1.3\text{GHz}}^{3\text{GHz}}$) to measure FIRRC and obtain a median $q_{\text{IR}} = 2.46_{-0.02}^{+0.03}$ and a standard deviation of $\sigma = 0.19$. In addition, we assume a fixed radio spectral slope with $\alpha = -0.8$ at rest-frame 0.4–10 GHz to k -correct the observed flux density at 1.3 GHz and obtain a median $q_{\text{IR}} = 2.45_{-0.04}^{+0.03}$ and a scatter of $\sigma = 0.18$. Therefore, for these SFGs

with $S_{3\text{GHz}} \geq 50 \mu\text{Jy}$ and $\text{SNR} > 3$ at GMRT 325 MHz and MeerKAT 1.3 GHz, the adoption of a fixed or an individually measured radio spectral index at high frequency does not significantly affect the median q_{IR} . However, as described in Section 4.2.1, the cut of $\text{SNR} > 3$ at GMRT 325 MHz biases our sample towards the SFGs with a relatively steeper radio spectrum at observer-frame 0.33–3 GHz. This might be the reason that our median q_{IR} is slightly lower than the measurements of some previous works (e.g. Bell 2003; Ivison et al. 2010; Thomson et al. 2014; Molnár et al. 2021).

We show the q_{IR} based on radio spectral indices at low and high frequencies for these 166 SFGs in Fig. 5. For the 89 SFGs with $\text{SNR} > 3$ at both GMRT 325 MHz and MeerKAT 1.3 GHz, we find that 28 per cent (25/89) of them have the measured $(q_{325\text{MHz}}^{1.3\text{GHz}} - q_{1.3\text{GHz}}^{3\text{GHz}})$ larger than their combined uncertainties. Here, the ‘combined uncertainty’ is defined in an analogous way to the σ_{combined} of equation (2), using the uncertainties of $q_{325\text{MHz}}^{1.3\text{GHz}}$ and $q_{1.3\text{GHz}}^{3\text{GHz}}$. In addition, three SFGs have the estimated lower limit of $(q_{325\text{MHz}}^{1.3\text{GHz}} - q_{1.3\text{GHz}}^{3\text{GHz}})$ larger than their combined uncertainties. Therefore, for the SFGs brighter than $50 \mu\text{Jy}$ at 3 GHz, > 17 per cent (28/166) of their q_{IR} are significantly underestimated if we adopt the radio spectral index at high frequency in the k -correction.

4.3.2 FIR-radio correlation for all of the 3 GHz-selected SFGs

Unfortunately, the majority of the 3 GHz-selected SFGs in our sample lack the measured flux density at GMRT 325 MHz because of the low sensitivity of 325 MHz data. We thus first use $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ instead of $\alpha_{325\text{MHz}}^{3\text{GHz}}$ to k -correct the observed luminosity at MeerKAT 1.3 GHz and measure the FIRRC for the 2012 VLA 3 GHz-selected SFGs (Fig. 6). Among them, 161 have $1 < \text{SNR} \leq 3$ at MeerKAT 1.3 GHz. We therefore use the lower limit of their $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ to estimate the lower limit of their q_{IR} as shown in Fig. 6. In addition, for the 2012 radio-selected SFGs, 1835 of them have $\text{SNR}_{\text{IR}} > 3$. The SNR_{IR} is the combined SNR of ‘superdeblended’ flux densities from 100 μm -to-1.2 mm, which reflects the quality of the FIR to millimeter SEDs that are used in estimating the total infrared luminosity. We therefore only keep the 1835 SFGs with $\text{SNR}_{\text{IR}} > 3$ in the following analyses. Among them, 1717 also have $\text{SNR} > 3$ at MeerKAT 1.3 GHz. Their median $q_{\text{IR}} = 2.47 \pm 0.01$, with a standard deviation of $\sigma = 0.23$, is consistent with that of SFGs brighter than $50 \mu\text{Jy}$ at 3 GHz (Section 4.3.1).

As shown in the top panels of Fig. 6, there is a clear trend between q_{IR} and redshift/stellar mass of galaxies, namely, that q_{IR} slightly declines with redshift and stellar mass. The evolution of q_{IR} with redshift and/or stellar mass has been widely discussed in literature (e.g. Jarvis et al. 2010; Magnelli et al. 2015; Delhaize et al. 2017; Algera et al. 2020; Delvecchio et al. 2021). However, a complete study of the evolution of q_{IR} and its physical factors are beyond the scope of this work. As described in Section 3.2, we remove radio excess sources from our sample to select a clean sample of SFG, which affects the completeness of study FIRRC of radio-selected sources. In this work, we only focus on how different k -corrections affect the study of FIRRC of SFGs.

To compare the FIRRC based on radio spectral index at low and high frequencies for all the 3 GHz-selected SFGs, we first divide our sample into the same redshift and stellar mass bins as those listed in Table 1 and show the median $q_{1.3\text{GHz}}^{3\text{GHz}}$ of the galaxies within each bin in Fig. 6. We also adopt the radio spectral index at low frequency, $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$, from our stacking analyses (Table 1) to k -correct the observed flux density at MeerKAT 1.3 GHz of galaxies within each bin and show the median $q_{325\text{MHz}}^{1.3\text{GHz}}$ in Fig. 6. We first perform

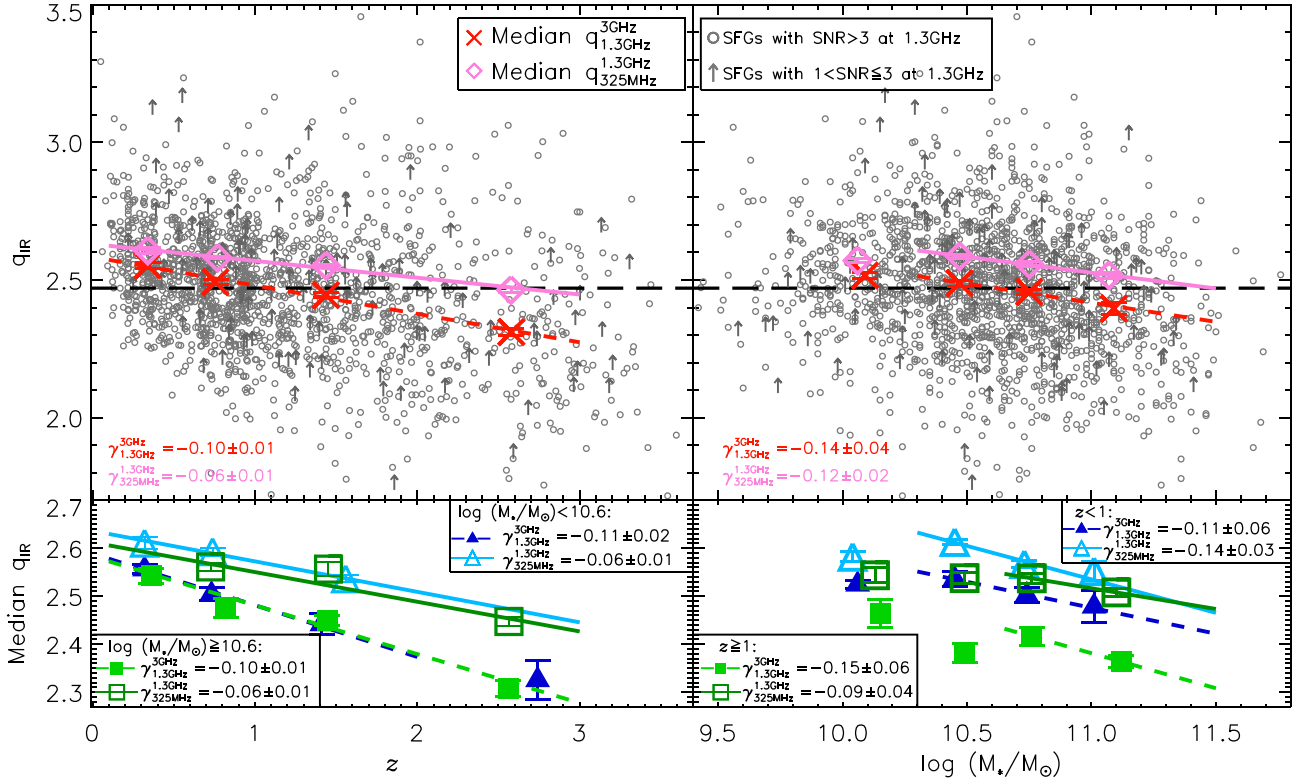


Figure 6. *Top:* Logarithmic infrared-to-radio luminosity ratio, q_{IR} as functions of redshift and stellar mass for the 1835 VLA 3 GHz-selected SFGs with measured spectral index $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ and $\text{SNR}_{\text{IR}} > 3$. Among them, 1717 also have $\text{SNR} > 3$ at MeerKAT 1.3 GHz. Their q_{IR} are measured by adopting $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ in k -correction. The black dashed line represents the median q_{IR} ($q_{\text{IR}} = 2.48 \pm 0.01$) of these SFGs. The red crosses show the median q_{IR} based on $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ of galaxies within each redshift and stellar mass bin while the error bars show the errors that are estimated from bootstrap resampling. For the remaining 118 SFGs with $1 < \text{SNR} \leq 3$ at MeerKAT 1.3 GHz, we use the lower limit of their $\alpha_{1.3\text{GHz}}^{3\text{GHz}}$ to estimate the lower limit of their q_{IR} as shown by upward arrows. We also adopt the averaged radio spectral index at low frequency ($\alpha_{325\text{MHz}}^{1.3\text{GHz}}$) from our stacking analyses (Table 1) for galaxies within each bin and show the median $q_{325\text{MHz}}^{1.3\text{GHz}}$ as magenta diamond. The solid and short-dashed lines represent the linear fits of correlations between median q_{IR} and physical properties based on the two different k -corrections, respectively. The slopes of these linear fits (γ) are shown on the bottom-left corner of each panel. We do not include the lowest stellar mass bin in the fits to reduce the effect of incompleteness of our flux-limited sample. *Bottom:* We further divide our sample into massive ($\log(M_*/M_\odot) \geq 10.6$) and less-massive ($\log(M_*/M_\odot) < 10.6$), high-redshift ($z \geq 1$) and low-redshift ($z < 1$) subsamples and show their median q_{IR} as functions of redshift and stellar mass. The solid symbols represent the median $q_{1.3\text{GHz}}^{3\text{GHz}}$ while the open symbols show the median $q_{325\text{MHz}}^{1.3\text{GHz}}$, which is estimated by adopting the averaged $\alpha_{325\text{MHz}}^{1.3\text{GHz}}$ from our stacking analyses. The highest redshift bin of less massive subsample, the lowest stellar mass bin of low-redshift subset, and the two lowest stellar mass bins of high-redshift subsample are excluded in the linear fits to reduce the effect of incompleteness of our radio flux-limited sample. Overall, the two different k -corrections result in a similar overall trend between q_{IR} and redshift or stellar mass, although evolution of q_{IR} with cosmic time will be slightly overestimated if we adopt the radio spectral index from high frequency ($\alpha_{1.3\text{GHz}}^{3\text{GHz}}$) in k -correction.

the simple linear fit to the correlations between median q_{IR} and physical properties of galaxies within each bin and overplot the results in Fig. 6. To be consistent with previous studies, we also fit the evolution of median q_{IR} as a function of redshift with the form $q \propto (1+z)^\gamma$ (e.g. Ivison et al. 2010; Delvecchio et al. 2021) and obtain the best-fitting $\gamma = -0.10 \pm 0.01$ for $q_{1.3\text{GHz}}^{3\text{GHz}}$ and $\gamma = -0.05 \pm 0.01$ for $q_{325\text{MHz}}^{1.3\text{GHz}}$. We further divide our sample into massive ($\log(M_*/M_\odot) \geq 10.6$) and less massive ($\log(M_*/M_\odot) < 10.6$) subsamples and show the median $q_{1.3\text{GHz}}^{3\text{GHz}}$ and $q_{325\text{MHz}}^{1.3\text{GHz}}$ as a function of redshift in the bottom-left panel of Fig. 6. We exclude the highest redshift bin of the less massive galaxies in the fit to reduce the effect of incompleteness in our sample. Overall, Fig. 6 and our analyses show that although the two different k -corrections result in a similar overall trend between q_{IR} and redshift, the evolution of q_{IR} with cosmic time based on radio spectral index at low frequency is slightly weaker than that based on spectral index derived from high frequency.

The bottom-right plot of Fig. 6 shows the median q_{IR} as a function of stellar mass for SFGs at high redshift ($z \geq 1$) and low redshift (z

< 1), respectively. To reduce the effect from incompleteness of our radio flux-limited sample, we do not include the lowest stellar mass bin and the two lowest stellar mass bins in fitting the correlations for the low-redshift and high-redshift subsamples, respectively. As shown in Fig. 6, the trends between q_{IR} and stellar mass based on the two radio spectral indices are consistent within the uncertainties.

Overall, using the radio spectral index from high frequency or the fixed spectral index of $\alpha = -0.8$ to k -correct the observed radio flux density will underestimate the q_{IR} for galaxies with a flatter radio spectrum at low frequency, and thus further affect the slope of the evolution of q_{IR} with redshift.

5 DISCUSSION

In this work, we have presented the radio spectral properties observed at 0.33–1.3 GHz and 1.3–3 GHz, respectively, for the 2094 VLA 3 GHz-selected SFGs. Here, we discuss the possible physical mechanisms that determine the radio spectrum at low and high frequencies.

5.1 Physical mechanisms that affect the radio spectrum at high frequency

In Section 4.1, we show that, for the 2012 SFGs with $z \sim 0.01\text{--}3$, the median radio spectral index between the observer-frame frequencies of 1.3–3 GHz is $\alpha_{1.3\text{ GHz}}^{3\text{ GHz}} = -0.78 \pm 0.01$. In addition, the median radio spectral index does not change significantly at rest-frame frequencies of 1.3–10 GHz, which suggests that the radio spectrum of SFGs at this frequency range is dominated by the synchrotron emission.

Fig. 2 shows that, on average, the radio spectrum at observer-frame 1.3–3 GHz slightly steepens with increasing stellar mass. One possible explanation for this correlation is the ‘aging’ of the relativistic electron population. In the literature, the CR electron population losing energy with time, i.e. via inverse-Compton, synchrotron, ionization, and bremsstrahlung process, is widely used to explain a steepening of the radio spectrum at high frequency (e.g. Condon 1992; Basu et al. 2015; Chyży et al. 2018; Klein et al. 2018; Thomson et al. 2019). However, for SFGs, the constantly injected CR electrons in star-forming regions will keep the radio spectral slope constant. Therefore, the slope of radio spectrum relies on the balance between the freshly injected and aged CR electrons in SFGs (e.g. Basu et al. 2015; Tabatabaei et al. 2017). The correlation between radio spectral index and stellar mass shown in Figs 2 and 4 could be explained by slightly increased ratio of aged and young relativistic CR electrons with increasing stellar mass of SFGs, although the data and the analyses in this work are not sufficient for us to be completely confident in this scenario.

We also find a weak positive correlation between sSFR and the radio spectral index at observer-frame 1.3–3 GHz of SFGs. We notice that the radio spectrum flattening with increasing sSFR has also been observed in local ULIRGs (Condon et al. 1991; Murphy et al. 2013). However, our results shown in Fig. 2 suggest that this trend in our sample might be an alternative manifestation of the underlying physics as indicated by the correlation between the radio spectrum and the stellar mass of SFGs. Murphy et al. (2013) suggested that galaxies with high sSFR are more compact and host deeply embedded star formation, thus more optically thick in the radio. Therefore, the radio spectrum flattening with increasing sSFR could be explained by increased free–free absorption (Murphy et al. 2013). However, further studies are necessary to explore the underlying physics that drive the correlations between radio spectral index and stellar mass, and sSFR of SFGs.

5.2 Physical mechanisms that flatten the radio spectrum at low frequency

Our analyses for all of the 3 GHz-selected SFGs and the subset brighter than 50 μJy at 3 GHz both show that, on average, the radio spectrum of SFGs flattens at low frequency ($\nu_{\text{obs}} < 1.3\text{ GHz}$). Some recent works on the radio spectral properties of SFGs have also reported a flatter radio spectrum at low frequency (e.g. Calistro Rivera et al. 2017; Thomson et al. 2019; Tisanić et al. 2019), although their sample selections and/or observer-frame frequencies are different from this work.

There are several physical mechanisms that might explain a flattening of the radio spectrum at low frequency, such as thermal absorption, the intrinsic curvature in the synchrotron spectrum due to energy losses of CR electrons and propagation effects, Razin effect, relativistic bremsstrahlung, synchrotron self-absorption, and so on (e.g. Condon 1992; McDonald et al. 2002; Tingay 2004; Murphy 2009; Clemens et al. 2010; Lacki 2013; Marvil et al. 2015; Kapińska

et al. 2017; Chyży et al. 2018; Klein et al. 2018). Among them, the thermal absorption and the intrinsic curvature in the synchrotron spectrum are the two main mechanisms that are widely included in the theoretical models in the literature. Models of the thermal free–free absorption assume that a power-law synchrotron radio spectrum is attenuated by free–free absorption with an absorption coefficient:

$$\left(\frac{\kappa}{\text{pc}^{-1}}\right) \approx 3.3 \times 10^{-7} \left(\frac{n_e}{\text{cm}^{-3}}\right)^2 \left(\frac{T_e}{10^4 \text{ K}}\right)^{-1.35} \left(\frac{\nu}{\text{GHz}}\right)^{-2.1}, \quad (4)$$

where n_e is the free electron density, T_e is the electron temperature, and ν is the radio frequency (Condon 1992).

Alternatively, models based on the intrinsic curvature in the synchrotron spectrum usually assume that the shock-accelerated CRs in the supernova remnants (SNRs) have a power-law energy distribution with spectral index $\alpha \sim -0.5$, which is supported by the averaged radio spectral index at $\sim 1\text{ GHz}$ of SNRs (e.g. Green 2014; Marvil et al. 2015; Chyży et al. 2018; Klein et al. 2018). As the relativistic CR electrons propagate within the galaxy, two major mechanisms cause energy loss at high frequency, namely synchrotron and inverse-Compton losses, which steepens the radio spectrum of SFGs. Although the remaining mechanisms may also flatten the radio spectrum at low frequency, previous work suggested that their impact in SFGs are very weak at rest-frame $\sim 0.4\text{--}1\text{ GHz}$ (e.g. Tingay 2004; Lacki 2013; Chyży et al. 2018).

In our work, as shown in Figs 2 and 4 and Table 1, on average, the radio spectrum at both low and high frequency slightly steepens with increasing stellar mass, which could be explained by age-related synchrotron losses. These results lend support to the models that take into account energy-loss mechanisms in explaining the flattening of the radio spectrum at low frequency. However, some recent studies based on nearby bright galaxies suggest that the different regions within a galaxy might have distinctly different radio spectra because of the inhomogeneous physical conditions within a galaxy (e.g. Marvil et al. 2015; Kapińska et al. 2017). Therefore, models with a single mechanism may not be adequate for interpreting the integrated radio spectrum of a galaxy.

5.3 Limitations

As described in Section 3, using 3 GHz as the selection frequency in this work provides a less-biased sample for studying radio spectrum at observer-frame 1.3–3 GHz. However, there might be some 1.3 GHz sources with steep radio spectra, which would not have been included in this work because of the non-detection at 3 GHz.

For studying the radio spectrum at observer-frame 0.33–3 GHz, the main limitation is the small sample size of SFGs with measured radio spectral index at both low and high frequency because of the low sensitivity of GMRT 325 MHz data. Although we applied a higher SNR cut and an additional flux density cut at 3 GHz to provide a less-biased sample when studying the radio spectral properties of SFGs, nearly half (46 per cent) of SFGs with a measured radio spectral index at low frequency ($\alpha_{325\text{ MHz}}^{1.3\text{ GHz}}$) have an SNR < 3 at 325 MHz. This leads to large uncertainties on the measured radio spectral index at low frequency and reduces the accuracy of the follow-up analyses. On the other hand, the cut of SNR > 3 at 325 MHz bias our sample towards the SFGs with relatively steeper radio spectra at observer-frame 0.33–3 GHz.

The other limitation is the uncertainties of the ‘superdeblended’ flux densities extracted from imaging data with very different angular resolutions. As shown in Fig. 1, although the ‘superdeblended’ flux densities extracted from VLA 1.4 GHz and MeerKAT 1.3 GHz imaging data are consistent with each other within the uncertainties,

the scatter is large. To improve quality and sample size in studying the radio spectral properties at both low and high frequency, observations with similar spatial resolutions and high sensitivity at multiple frequencies are essential. The recently released high-sensitivity low-frequency (<150 MHz) data from the Low Frequency Array (LOFAR) large surveys (e.g. de Gasperin et al. 2021; Sabater et al. 2021; Tasse et al. 2021) and the undergoing joint project, SuperMIGHTEE, which combines the MeerKAT and uGMRT telescopes to observe the MIGHTEE survey regions at 0.5–2.7 GHz with the same angular resolution (~ 5 arcsec) and similar sensitivity, will dramatically advance the studies of the radio spectrum.

6 CONCLUSION

By combining deep MeerKAT 1.3 GHz data with VLA 3 GHz, GMRT 325 MHz data and rich ancillary data available in the COSMOS field, we study the radio spectral properties of 2094 radio 3 GHz-selected SFGs within the MIGHTEE-COSMOS early science coverage. Our main conclusions are as follows.

(i) After removing AGNs and red quiescent galaxies from the 4014 VLA 3 GHz-selected sources within the MIGHTEE-COSMOS early science coverage, we obtain 2094 SFGs with $\text{SNR} > 5$ at 3 GHz. We extract flux densities from the MeerKAT imaging data for these 3 GHz-selected SFGs and obtain the radio spectral indices between the observer-frame frequencies of 1.3–3 GHz for 96 per cent (2012/2094) of them. We limit our sample to the SFGs with $\text{SNR} > 3$ at MeerKAT 1.3 GHz and obtain a median spectral index of $\alpha_{1.3\text{ GHz}}^{3\text{ GHz}} = -0.80 \pm 0.01$. Although the scatter is large, the median radio spectral slope of these SFGs does not change within rest-frame ~ 1.3 –10 GHz. This confirms that on average the radio spectrum of SFGs at rest-frame ~ 1.3 –10 GHz is dominated by synchrotron emission.

(ii) On average, the radio spectrum at observer-frame 1.3–3 GHz slightly steepens with increasing stellar mass, and flattens with sSFR of SFGs, which may reflect the same fundamental relation. These trends could be explained by age-related synchrotron losses, i.e. the CR electron population losing energy with time, steepening the radio spectrum at rest-frame ~ 1.3 –10 GHz of SFGs.

(iii) We extract flux densities from the GMRT 325 MHz data for the 3 GHz-selected SFGs in our sample. Due to the lower sensitivity of GMRT 325 MHz data, we apply a further flux density cut ($S_{325\text{ MHz}} \geq 50 \mu\text{Jy}$) to obtain a less-biased sample to study the radio spectral properties at observer-frame 0.33–3 GHz. The median radio spectral indices for the 166 SFGs with measured flux densities in the three radio frequencies are $\alpha_{325\text{ MHz}}^{1.3\text{ GHz}} = -0.59^{+0.02}_{-0.03}$ and $\alpha_{1.3\text{ GHz}}^{3\text{ GHz}} = -0.74^{+0.01}_{-0.02}$. Therefore, on average, the radio spectrum of SFGs flattens at low frequency (0.33–1.3 GHz).

(iv) For the SFGs that lack measured flux densities at GMRT 325 MHz, we stack both GMRT 325 MHz and MeerKAT 1.3 GHz imaging data at their 3 GHz positions and estimate their averaged radio spectral index at observer-frame 0.33–1.3 GHz within each redshift, stellar mass, SFR, sSFR, and high-frequency radio spectral index ($\alpha_{1.3\text{ GHz}}^{3\text{ GHz}}$) bins. Our stacking analyses show that the averaged radio spectral index also steepens with increasing stellar mass and flattens with sSFR of SFGs. In addition, the radio spectral indices at low and high frequencies have a positive correlation.

(v) We use the two different spectral indices measured at low (0.33–1.3 GHz) and high frequencies (1.3–3 GHz) to k -correct the observed flux density at MeerKAT 1.3 GHz, and thus obtain two sets of q_{IR} values for the SFGs in our sample. For the SFGs with measured flux density at the three radio frequencies, using the high-frequency spectral index for k -correction underestimates the q_{IR} for

>17 per cent of them that have a flatter radio spectrum at low frequency. For all of the 3 GHz-selected SFGs, using the spectral index at high frequency in k -correction slightly overestimate the evolution of q_{IR} with cosmic time, although both choices of the radio spectral index (at low or high frequency) result in similar trends for q_{IR} -redshift and q_{IR} -stellar mass relations. Therefore, deep low-frequency radio continuum data are essential for an accurate study of FIRRC of SFGs.

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DATA AVAILABILITY

The MeerKAT 1.3 GHz data were accessed from the South African Radio Astronomy Observatory (SARAO, www.ska.ac.za). The derived data generated in this research will be shared on reasonable request to the corresponding author.

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